

# Vortex dynamics and irreversibility line in optimally doped $\text{SmFeAsO}_{0.8}\text{F}_{0.2}$ from ac susceptibility and magnetization measurements

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Ac susceptibility and static magnetization measurements were performed in the optimally doped  $\text{SmFeAsO}_{0.8}\text{F}_{0.2}$  superconductor. The field - temperature phase diagram of the superconducting state was drawn and, in particular, the features of the flux line lattice derived. The dependence of the intra-grain depinning energy on the magnetic field intensity was derived in the thermally-activated flux creep framework, enlightening a typical  $1/H$  dependence in the high-field regime. The intra-grain critical current density was extrapolated in the zero temperature and zero magnetic field limit, showing a remarkably high value  $J_{c0}(0) \sim 2 \cdot 10^7 \text{ A/cm}^2$ , which demonstrates that this material is rather interesting for the potential future technological applications.

PACS numbers: 74.25.Uv, 74.25.Wx, 74.70.Xa

## I. INTRODUCTION

Several properties of iron-based 1111 oxy-pnictide superconductors like, for instance, the high crystallographic and superconducting anisotropy and large penetration depths make them similar to cuprate high- $T_c$  materials. Since their discovery,<sup>1,2</sup> several studies have been performed in order to clarify their intrinsic microscopic properties.<sup>3,4</sup> The attention has been mainly focussed on the bosonic coupling mechanism of the superconducting electrons, on the features of the spin density wave magnetic phase characterizing the parental and lightly doped-compounds and its possible coexistence with superconductivity,<sup>5-7</sup> and on the interaction between localized magnetism from RE ions and itinerant electrons onto FeAs bands.<sup>8-10</sup> A common hope is that a full comprehension of these topics in oxy-pnictide superconductors could also allow one to answer several open questions on the cuprates.

At the same time, other analogies with cuprates possibly characterize 1111 oxy-pnictides as interesting materials for technological applications, like small coherence lengths (and, correspondingly, high values of upper critical fields) besides the high values of the superconducting critical temperature  $T_c$ . In this respect, studies of macroscopic properties like critical fields and critical depinning currents are of the utmost importance. Namely, the investigation of the dynamical features of the flux line lattice and of the so-called irreversibility line, typically investigated by means of both resistance and ac susceptibility measurements, is in order. Those measurements allow to further check the validity of the theories used to model the mixed state of cuprate materials and, in particular, the vortices motion and its relationship with the possible pinning mechanisms.<sup>11</sup> Several works reporting magnetoresistance,<sup>12-14</sup> modulated microwave absorption<sup>15</sup> and dc magnetization<sup>16-19</sup> examining the flux line lattice dynamics in 1111 oxy-pnictides have al-

ready been published in the last two years. To our knowledge, no study of the flux line lattice by means of ac susceptibility measurements has been published yet.

This paper deals with the field, temperature and frequency dependences of ac susceptibility in optimally doped  $\text{SmFeAsO}_{0.8}\text{F}_{0.2}$ , which is one of the compounds showing the highest  $T_c$  among all the iron-based superconductors. Although no large enough single-crystals are available and our data refer to unoriented powder samples, the power of the ac susceptibility technique allowed us to deduce several intrinsic features of the mixed state of the superconductor. The magnetic field ( $H$ ) behaviour of the irreversibility line was obtained, allowing to draw (together with dc magnetization data) a detailed phase diagram of the flux line lattice. The  $H$  dependence of the intra-grain effective depinning energy ( $U_0$ ) was investigated, evidencing the characteristic crossover from a single-vortex-dynamics to a collective-dynamics ( $U_0 \propto 1/H$ ) at a field  $H \sim 0.5 \text{ T}$ . An estimate of the intra-grain critical depinning current density in the limit of vanishing temperature and magnetic field was also deduced, giving the remarkably high value of  $J_{c0}(0) \sim 2 \cdot 10^7 \text{ A/cm}^2$ . This result is of great importance in characterizing  $\text{SmFeAsO}_{0.8}\text{F}_{0.2}$  as a superconductor suitable for technological applications.

## II. EXPERIMENTALS AND MAIN RESULTS

$\text{SmFeAsO}_{0.8}\text{F}_{0.2}$  was prepared by solid state reaction at ambient pressure from Sm, As, Fe,  $\text{Fe}_2\text{O}_3$  and  $\text{FeF}_2$ . SmAs was first synthesized from pure elements in an evacuated, sealed glass tube at a maximum temperature of  $550^\circ\text{C}$ . The final sample was synthesized by mixing SmAs, Fe,  $\text{Fe}_2\text{O}_3$  and  $\text{FeF}_2$  powders in stoichiometric proportions, using uniaxial pressing to make powders into a pellet and then heat treating the pellet in an evacuated, sealed quartz tube at  $1000^\circ\text{C}$  for 24 hours, followed by furnace cooling. The sample was analyzed by powder

X-ray diffraction in a Guinier camera, with Si as internal standard. The powder pattern showed the sample to be single phase with two weak extra lines at low angle of the SmOF extra phase. The lattice parameters were  $a = 3.930(1)$  Å and  $c = 8.468(2)$  Å.

Static magnetization  $M_{dc}$  measurements were performed by means of a Quantum Design MPMS-XL7 SQUID magnetometer. The temperature  $T$  dependence of  $M_{dc}$  upon field-cooling (FC) the sample was monitored at different applied magnetic fields up to 7 T. Representative raw  $M_{dc}$  vs.  $T$  curves are shown in Fig. 1. The su-

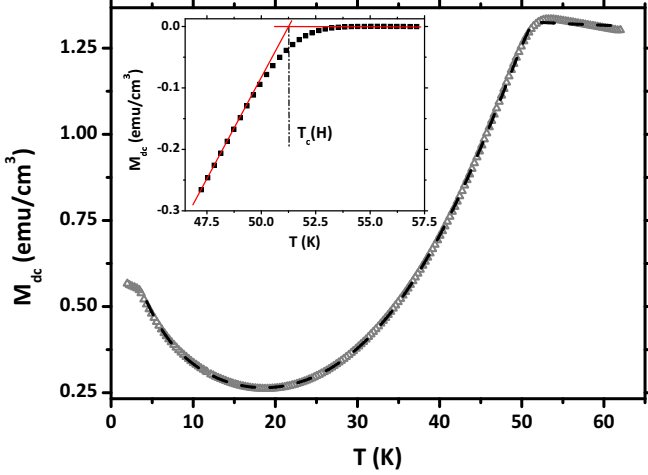


FIG. 1: (Color online) Temperature dependence of static magnetization  $M_{dc}$  upon field-cooling the sample at  $H = 1.5$  T. Dashed lines are the best fitting functions according to Eq. 2 (see later on in the text). Inset: estimate of  $T_c(H = 1.5$  T) after the subtraction of the paramagnetic contribution.

perconducting (SC) response, with onset around  $T_c \simeq 52$  K, is found to be superimposed to a paramagnetic contribution associated with  $\text{Sm}^{3+}$  ions. Clear kinks in the magnetization curves can be observed at  $T_N \simeq 4$  K, evidencing the antiferromagnetic transition of the  $\text{Sm}^{3+}$  magnetic moments.<sup>20</sup> The field dependence of the SC transition temperature  $T_c(H)$  was deduced by first subtracting the linear extrapolation of the  $\text{Sm}^{3+}$  paramagnetic contribution in a few-K region around the SC onset from the raw data. The transition temperatures were then estimated from the intersection of two linear fits of the resulting curves above and below the onset (see Fig. 1, inset).  $T_c(H)$  behaviour was also deduced by means of magnetoresistance measurements upon the application of magnetic fields up to 9 T, showing a behaviour analogous to what observed in  $\text{SmFeAsO}_{1-x}\text{F}_x$  compounds from the same bath but with lower  $x$  concentrations of  $\text{F}^-$ .<sup>21</sup>

The onset of the diamagnetic contribution and its dependence on the applied external magnetic field was also investigated by means of a Quantum Design MPMS-XL5 SQUID ac susceptometer. Measurements were performed with an alternating field in the range  $H_{ac} = (0.0675 - 1.5) \cdot 10^{-4}$  T parallel to the static field  $H$ , which varied up to

5 T. The ac field frequency ranged from 37 to 1488 Hz. The diamagnetic onset temperature was estimated from  $\chi'$  vs.  $T$  curves by means of the same procedure shown in the inset of Fig. 1.

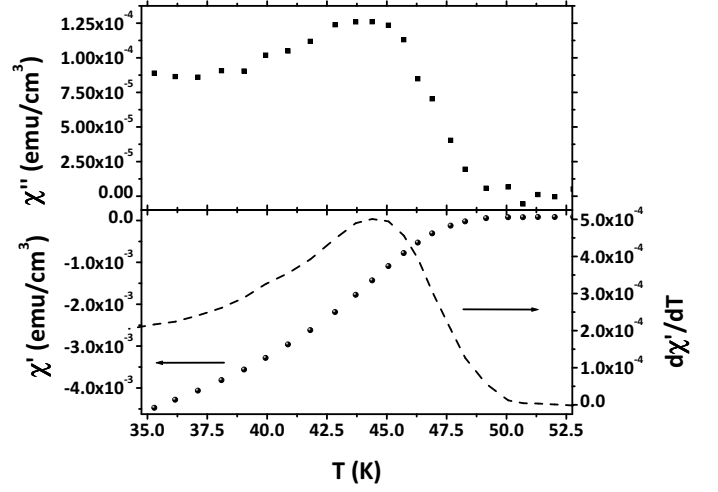


FIG. 2: Temperature dependence of the real and imaginary components of ac susceptibility  $\chi'$  and  $\chi''$  (bottom and top panels, respectively). The alternating field amplitude  $H_{ac} = 1.5 \cdot 10^{-4}$  T at the frequency  $\nu_m = 37$  Hz is superimposed to a larger static field  $H = 1.5$  T. The dashed line in the bottom panel represents the first derivative of  $\chi'$  with respect to temperature.

An accurate examination of ac susceptibility data as a function of  $\nu_m$  allowed us to obtain further information on the dynamical properties of the flux line lattice (FLL). It is well known, in fact, that a peak in  $\chi''$  vs.  $T$  curves associated with a maximum in the energy dissipation inside the sample appears at a temperature  $T_p$  slightly lower than the diamagnetic onset temperature in  $\chi'$ . Correspondingly, at the same temperature  $T_p$  the  $\chi'$  vs.  $T$  curve displays a maximum in its first derivative<sup>22</sup> (see exemplifying raw data in Fig. 2, lower panel). Several works in the past decades have tried to clarify the origin of the  $\chi''$  peak. One of the possible interpretations relies on the Bean's critical state model<sup>23,24</sup> and associate the peak in  $\chi''$  with the flux front reaching the centre of the sample. In this case, the peak temperature  $T_p$  should not depend on the measurement frequency  $\nu_m$  and a strong dependence on sample dimensions and ac field amplitude  $H_{ac}$  are predicted.<sup>25</sup> Another interpretation relies on the modification of the skin depth, due to the superconductor resistivity in the thermally-assisted flux flow (TAFF) regime, with respect to the London penetration depth. In this case  $T_p$  should strongly depend on the measurement frequency  $\nu_m$  while no dependence on the ac field amplitude  $H_{ac}$  is predicted.<sup>22,25-29</sup> Considering the frequency dependence of the  $\chi''$  peak, another interesting interpretation has been associated with a resonant absorption of energy obtained when the inverse of  $\nu_m$  matches the characteristic relaxation time  $\tau_c$  of the

FLL at  $T_p$ ,<sup>30,31</sup> namely

$$2\pi\nu_m\tau_c|_{T=T_p} = 1. \quad (1)$$

In this case, the underlying theory is the more general framework of the thermally activated creep of flux lines between different metastable minima of pinning potentials.<sup>32,33</sup>

At temperatures lower than  $T_p$ , other broader contributions to both  $\chi'$  and  $\chi''$  were found and interpreted as arising from granularity of the powder sample and, in particular, to intergranular Josephson weak links between different grains.<sup>29,34,35</sup> In cuprate materials, from the analysis of the low-temperature peak in  $\chi''$  and, in particular, of its frequency dependence, the depinning energy barrier associated with grain boundaries was extracted.<sup>34,36,37</sup> Strong granularity has been observed also in iron-based pnictide materials. On samples prepared with the same procedure a small (though not negligible) intergranular critical density current has been evaluated by a remanent magnetization analysis.<sup>38</sup> However, it has been determined that the main contribution to the magnetization curve comes from intragranular currents. By considering our data just in the  $T$  region close enough to the diamagnetic onset, then, we will be focussing only on the intra-grain intrinsic dynamical properties. The stronger signal amplitude, moreover, made it reasonable to investigate the peak in  $\chi'$  derivative instead of the maximum in  $\chi''$ .

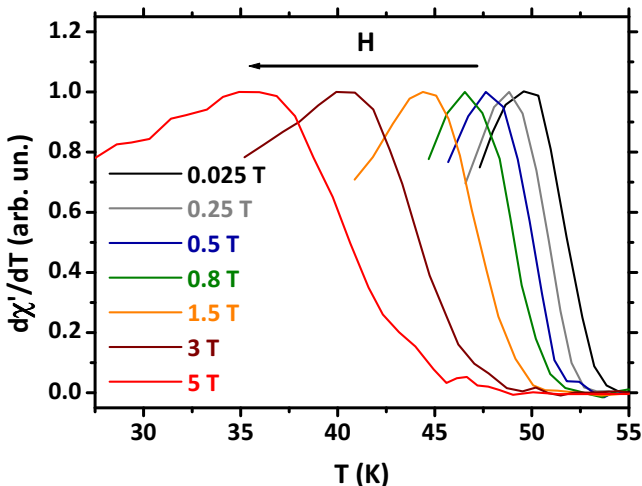


FIG. 3: (Color online) Temperature dependence of the first derivative of the real component of ac susceptibility  $\chi'$  ( $H_{ac} = 1.5 \cdot 10^{-4}$  T,  $\nu_m = 37$  Hz) with respect to temperature at different applied magnetic fields. All the single curves were independently normalized to the corresponding maximum derivative value.

Fig. 3 shows the temperature evolution of the normalized first derivative of the real component  $\chi'$  of the ac susceptibility upon the application of different values of static magnetic field  $H$ . In these measurements, both the alternating field and the working frequency were kept

fixed at  $H_{ac} = 1.5 \cdot 10^{-4}$  T and  $\nu_m = 37$  Hz respectively. The effect of increasing  $H$  is a clear shift of  $T_p$  towards lower values. At each applied  $H$  a clear dependence of  $T_p$  on  $\nu_m$  was evidenced, as discussed later on. Some scans with  $H_{ac}$  values in the range  $(0.0675 - 1.5) \cdot 10^{-4}$  T were also performed at the representative values  $H = 0.025, 0.25$  and 5 T (data not shown). Within the experimental error, no dependence of  $T_p$  on  $H_{ac}$  was evidenced.

### III. ANALYSIS AND DISCUSSION

The raw  $M_{dc}$  vs.  $T$  data reported in Fig. 1 were fitted by the function (see dashed curves in Fig. 1)

$$M_{dc}(T, H) = M_{sc} \left[ 1 - \left( \frac{T}{T_c} \right)^\alpha \right]^\beta + C_{cw} \frac{H}{T - T_N} + M_0(H) \quad (2)$$

where the first term is the diamagnetic Meissner response (empirically represented by a two-exponents mean-field function) and the second one is the Curie-Weiss paramagnetic contribution. The last term accounts for all the sources of  $T$ -independent magnetism, ranging from Pauli- and Van-Vleck-like paramagnetism to a small contribution of magnetic impurities (e.g. Fe<sub>2</sub>As).<sup>39</sup> A detailed analysis of the results of the fitting procedure according to Eq. 2 will be presented in another work.<sup>40</sup>

Results from both SQUID magnetometry and ac susceptibility are summarized in the phase diagram shown in Fig. 4.

From the  $T_c(H)$  data obtained from  $M_{dc}$  vs.  $T$  curves (see Fig. 1) it was possible to derive the temperature dependence of the upper critical field  $H_{c2}$  (see open circles in Fig. 4). A linear fit of the  $H_{c2}$  vs.  $T$  data deduced from magnetization data gives a slope  $dH_{c2}/dT = 7.47 \pm 0.15$  T/K, in agreement with what was found in compounds of the same family from magnetoresistivity data,<sup>21,41</sup> even if much lower slope values were reported from calorimetric measurements on single crystals of Nd-based 1111 superconductors.<sup>42</sup> Then, in the simplified assumption of single-band superconductivity, through the Werthamer, Helfand and Hohenberg relation<sup>43</sup>

$$H_{c2}(T = 0K) \simeq 0,693 \cdot T_c \left| \frac{dH_{c2}}{dT} \right|_{T \simeq T_c} \quad (3)$$

it is possible to estimate  $H_{c2} \simeq 270$  T in the limit of vanishing  $T$ .

A comparison between the diamagnetic onset temperature as obtained from ac susceptibility at different frequencies (see exemplifying raw data in Fig. 2) and the one obtained from static magnetization is plotted in the inset of Fig. 4. The onset in ac data was systematically found at lower temperatures than the corresponding dc diamagnetic onset. Dashed lines represent the empirical

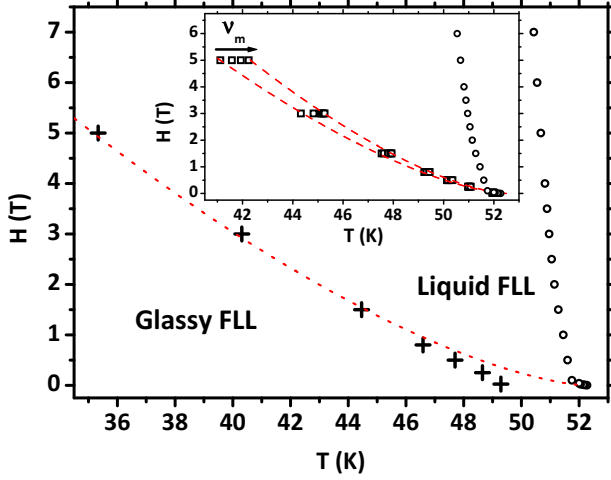


FIG. 4: (Color online) Phase diagram associated with the flux line lattice. A glassy phase of FLL is separated from a liquid phase by the irreversibility line obtained from the field dependence of the maximum slope of  $\chi'$  (plus signs (+), relative to  $H_{ac} = 1.5 \cdot 10^{-4}$  T,  $\nu_m = 37$  Hz). The red dotted line is a best fit according to Eq. 4. Circles (o) track the behaviour of  $H_{c2}$  as measured from dc magnetization data. Inset: comparison of the diamagnetic onset measured from dc magnetization (o) and ac susceptibility ( $\square$ ) with  $H_{ac} = 1.5 \cdot 10^{-4}$  T and  $\nu_m = 37 - 1488$  Hz. The red dashed lines are best fits of data at  $\nu_m = 37$  Hz and  $\nu_m = 1488$  Hz according to Eq. 4.

power-law fitting function

$$H \propto (1 - T/T_c)^\beta, \beta = 3/2 \quad (4)$$

well describing experimental data at each value of the ac field frequency  $\nu_m$ . Such a functional form characterized by  $\beta = \frac{4}{3} - \frac{3}{2}$  is known to describe the irreversibility line in the  $H - T$  phase diagram of cuprates.<sup>44</sup>

Plus signs in the phase diagram in Fig. 4 refer to the points  $T = T_p$  of maximum slope of  $\chi'$  vs.  $T$  curves ( $H_{ac} = 1.5 \cdot 10^{-4}$  T and  $\nu_m = 37$  Hz. See Fig. 3) corresponding, as already explained in sect. II, to the maximum of  $\chi''$  vs.  $T$  associated with intrinsic intra-grain losses. These data divide the phase diagram into two main regions, following the interpretation of the  $\chi''$  peak in term of resonant absorption of energy in a thermally activated flux creep model.<sup>30</sup> In the high- $T$  and high- $H$  region the flux lines are in a reversible state, that is, they are responding to the external ac perturbation (liquid FLL). On the other hand, in the low- $T$  and low- $H$  region vortices are arranged in a glassy-like frozen FLL that gives rise to a non-reversible response and to dissipation, linked to the non-zero values assumed by  $\chi''$ .  $T_p$  vs.  $H$  points associated with the lowest accessible frequency  $\nu_m$  are thus expected to belong to the irreversibility line (or de Almeida - Thouless line) of the FLL phase diagram. As in the case of the diamagnetic onset in  $\chi'$  (see Fig. 4, inset), Eq. 4 nicely fits the field dependence of  $T_p$  in the  $H > 0.8$  T limit (see the dotted line in Fig. 4).

A logarithmic dependence of  $1/T_p$  vs.  $\nu_m$  at every fixed  $H$  is evidenced over the explored frequency range (see, for example,  $H = 1.5$  T in the inset of Fig. 5). Data can then be fitted within a thermally-activated framework by the formula (red dashed line in the inset of Fig. 5)

$$\frac{\nu_m}{\nu_0} = \exp \left( -\frac{U_0(H)}{k_B T_{p|\nu=\nu_m}} \right) \quad (5)$$

from which it can be observed that the logarithmic behaviour of  $1/T_p$  is mainly controlled by the parameter  $U_0$ , playing the role of an effective depinning energy barrier in a thermally-activated flux creep model. The parameter  $\nu_0$  takes the meaning of a intra-valley characteristic frequency associated with the motion of the vortices around their equilibrium position in the pinning centers.

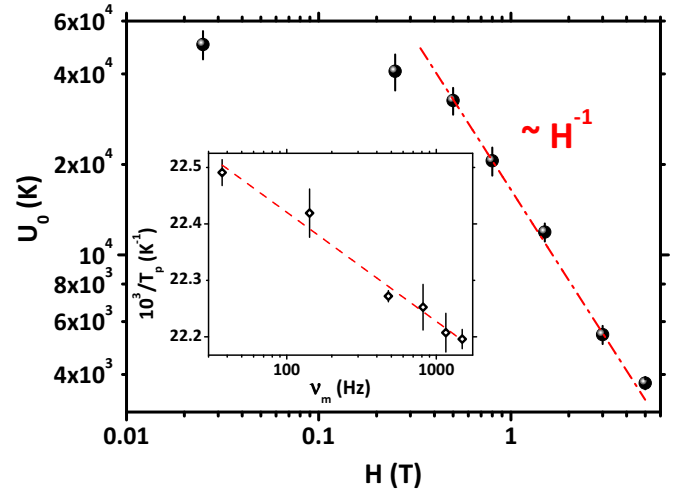


FIG. 5: (Color online) Magnetic field dependence of the effective depinning energy barrier. The red dashed-dotted line is a best fit of data according to a  $1/H$  dependence for  $H \gtrsim 0.5$  T. Inset: frequency dependence of  $1/T_p$ . Data refer to  $H = 1.5$  T with the alternating field amplitude fixed to  $H_{ac} = 1.5 \cdot 10^{-4}$  T. The red dashed line is a best fit of the experimental data according to Eq. 5.

The advantage of extracting the value of  $U_0$  from ac susceptibility data, if compared for instance to magnetoresistance data, is that the former is an almost isothermal estimate. Temperature, in fact, varies at most by 1 K degree as a function of  $\nu_m$  at the highest applied  $H$  (see Fig. 5, inset) so that it is possible to determine  $U_0$  at a temperature  $T^*(H)$  with a maximum uncertainty of 0.5 K. This fact will be of interest when deriving the critical current density value, as it will be shown later on.

Data at different static magnetic fields can be fitted according to this model giving the results reported in Fig. 5, where the  $H$  dependence of the effective depinning energy barrier  $U_0$  is shown. Beyond an overall sizeable reduction of  $U_0$  with increasing  $H$ , a crossover between two different regimes can be clearly observed at

$H \gtrsim 0.5$  T. At high fields the data are well described by a  $1/H$  dependence, a well-known result in high- $T_c$  cuprate superconductors, observed by means of several techniques, ranging from NMR<sup>45</sup> to ac magnetometry<sup>46</sup> and magnetoresistivity<sup>30</sup>. A naive explanation of this behaviour can be obtained in terms of the balance between the Gibbs free energy of the system and the energy required for the motion of a flux lines bundle.<sup>47</sup> In this framework, the crossover between the two different trends of  $U_0$  vs.  $H$  shown in Fig. 5 can be interpreted as the transition from a basically single-flux line response at low  $H$  values to a collective response of vortices for  $H > 0.5$  T. A similar phenomenology was recently reported from magnetoresistivity data on a single crystal of O-deficient  $\text{SmFeAsO}_{0.85}$  and on powder samples of La-based and Ce-based 1111 superconductors.<sup>13,14</sup> The crossover between different regimes was observed at  $H \simeq 1$  T in La-based samples and at much higher magnetic fields in Ce-based samples and in  $\text{SmFeAsO}_{0.85}$  ( $H \simeq 3$  T). However,  $U_0$  values are typically 1 order-of-magnitude lower in La- and Ce- based superconductors. Comparing the sets of data for  $U_0$  derived from magnetoresistivity and here from ac susceptibility in Sm-based samples, the numerical agreement is very good for  $H \gtrsim 1$  T.

In the  $1/H$  regime, Tinkham<sup>47</sup> extended previous works by Yeshurun and Malozemoff<sup>44,48</sup> showing that the relation for the normalized effective depinning energy barrier

$$\frac{U_0(t, H)}{t} \simeq \frac{K J_{c0}(0)}{H} g(t) \quad (6)$$

holds also for granular samples. Here  $U_0$  is expressed in K degrees,  $t$  is the reduced temperature  $t \equiv T/T_c$ ,  $g(t) \equiv 4(1-t)^{3/2}$ ,  $J_{c0}(t)$  quantifies the critical current density at  $H = 0$  and  $T = tT_c$ , while the constant  $K$  is defined as  $K \equiv 3\sqrt{3}\Phi_0^2\beta/2c$ ,  $\Phi_0$  being the flux quantum,  $c$  the speed of light and  $\beta$  a numerical constant close to unit value. Eq. 6 is derived in the simplified assumption of a two-fluid model.<sup>47</sup> By assuming that this empirical scenario can describe also the system under current investigation, from ac susceptibility it is possible to directly extrapolate the value of  $J_{c0}(0)$ .  $U_0(H)$ , as already noticed above, is almost isothermally estimated and can then be expressed as  $U_0(t^*, H)$ . By now plotting  $U_0/t^*g(t^*)$  as a function of  $1/H$  (see Fig. 6), from a linear fit of data it is possible to extract from Eq. 6 the value  $J_{c0}(0) = (2.25 \pm 0.05) \cdot 10^7$  A/cm<sup>2</sup>, having assumed  $\beta = 1$ . This rather high value is in agreement with estimates of the critical current density  $J_c$  evaluated by inductive measurements in similar Sm-based samples<sup>38</sup> and also with the direct measurement of this quantity in a Sm-based 1111 single crystals in the  $T \rightarrow 0$  K and  $H \rightarrow 0$  T limit.<sup>12</sup>

## IV. CONCLUSIONS

The  $H - T$  phase diagram of the flux line lattice in a powder sample of optimally-doped  $\text{SmFeAsO}_{0.8}\text{F}_{0.2}$  was

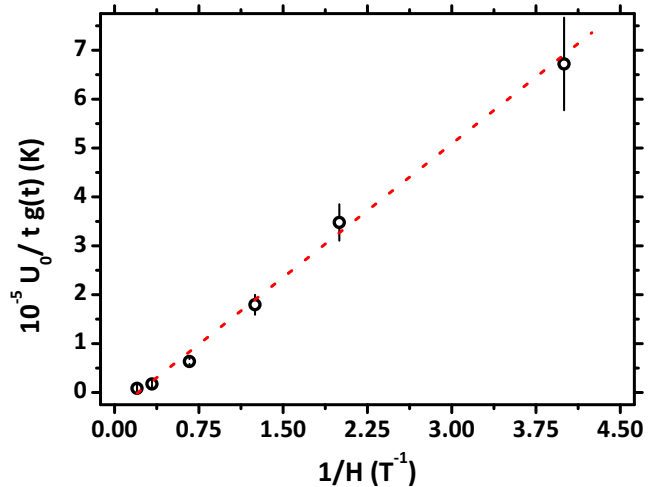


FIG. 6: (Color online) Plot of the effective energy barrier normalized with respect to the reduced temperature function  $tg(t)$  (see text) as a function of the inverse applied static field. The slope of the linear fit (see red dotted curve) is directly proportional to the critical current density extrapolated at  $T = 0$  K and  $H = 0$  T.

investigated by means of both ac and dc susceptibility measurements. The irreversibility line separating a liquid from a glassy phase was deduced and the activation depinning energy  $U_0$  as a function of the external magnetic field derived in the framework of a thermally-activated flux creep theory. A  $1/H$  dependence of  $U_0$  for  $H \gtrsim 0.5$  T, typical of collective motion of flux lines, has been evidenced. From the  $U_0$  vs.  $H$  behaviour a value of  $J_{c0}(0) \sim 2 \cdot 10^7$  A/cm<sup>2</sup> has been extrapolated for the critical depinning current at both zero field and zero temperature. From this result we confirm that high intrinsic critical depinning current density values seem to be a peculiar feature of these systems, possibly making them good candidates for technological applications.

## Acknowledgments

M. J. Graf, A. Rigamonti and L. Romanò are gratefully acknowledged for stimulating discussions. We thank C. Pernechele and D. Zola for useful help and suggestions about ac susceptibility measurements.

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